Renormalization Hopf Algebra Emerging from Stora's Main Theorem

Dimensional Regularization in Position Space and a Forest Formula for Epstein-Glaser Renormalization

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Motivation

A Brief History of Perturbative Renormalization Theory

Local Freedom in S-matrix Stückelberg & Petermann ('53): Popineau & Stora ('82): Main Theorem of Renormalization

Bogoliubov & Parasiuk ('57), Hepp ('66):

DimReg+MS

This Talk Epstein & Glaser ('73)

Kreimer & Connes ('98, '00)

algebra | Pinter ('00)

Hopf Y Gracia-Bondía & Lazzarini,

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Overview

- Perturbative Algebraic QFT (pAQFT)
- 2 Renormalization Group and Main Theorem
- 3 Renormalization Hopf algebra
- 4 DimReg+MS
- 5 Conclusion and Outlook

pAQFT - Algebra of Observables

Classical Field Theory

Observables in Classical Field Theory: $f,g \in \mathcal{D}(\mathbb{M})$

$$F: \mathscr{E}(\mathbb{M}) \to \mathbb{C}, \ \varphi \mapsto F(\varphi) \ \left[= \int f(x) \left(\partial \varphi(x) \right)^2 + g(x) \left(\varphi(x) \right)^4 dx \right]_{\text{e.g.}}$$

Smooth Functionals:

$$\forall n: \quad \left\langle F^{(n)}(\varphi), h^{\otimes n} \right\rangle = \frac{d^n}{d\lambda^n} F(\varphi + \lambda h) \Big|_{\lambda=0} \in \mathbb{C}, \quad h \in \mathscr{E}(\mathbb{M}).$$

Definition (Deformable Algebra)

Poisson algebra $(\mathcal{F}(\mathbb{M}), \lfloor \cdot, \cdot \rceil, \cdot)$, such that $\forall F, G \in \mathcal{F}(\mathbb{M})$:

- Pointwise Product: $(F \cdot G)(\varphi) := F(\varphi)G(\varphi)$.
- Poisson structure: $[F,G](\varphi) := \langle F^{(1)}, \Delta G^{(1)} \rangle (\varphi)$, where $\Delta = \Delta_{\mathsf{ret}} \Delta_{\mathsf{adv}}$.
- $\quad \circ \ F^{(n)}(\varphi) \in \mathscr{E}'(\mathbb{M}^n) \quad \text{and} \quad \left[\mathsf{WF}(F^{(n)}) \right]_{2^{\mathsf{nd}}} \cap \left\{ \overline{\mathcal{V}^+}^n \cup \overline{\mathcal{V}^-}^n \right\} = \emptyset.$

pAQFT - Algebra of Observables

Deformation Quantization

Bidifferential operator:

$$\Gamma_H := \int dx \ dy \ H(x,y) \, rac{\delta}{\delta arphi(x)} \otimes rac{\delta}{\delta arphi(y)} \, ,$$

 $H = \Delta_+ + H^{\text{sym}}$: Hadamard bisolution, $\Delta_+ = (i\Delta)^+$: Wightman func. H^{sym} : smooth, symmetric, Lorentz invariant Klein-Gordon bisolution.

Deformation:

$$\mathcal{F}(\mathbb{M})[[\hbar]]^{\otimes 2} \xrightarrow{\exp(\hbar\Gamma_{H})} \mathcal{F}(\mathbb{M})[[\hbar]]^{\otimes 2}, F \star G := \sum_{k=0}^{\infty} \frac{\hbar^{k}}{k!} \left\langle F^{(k)}, H^{\otimes k} G^{(k)} \right\rangle.$$

$$\mathcal{F}(\mathbb{M})[[\hbar]]$$

Algebra of Observables: $(\mathcal{F}(\mathbb{M})[[\hbar]], \star)$,

$$\frac{1}{\hbar}[F,G]_{\star} \xrightarrow{\hbar \to 0} [F,G]$$

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pAQFT - Time-Ordered Product

Second Order Differential Operator:

$$\Gamma'_{H_F} := \frac{1}{2} \int dx \, dy \, H_F(x,y) \, \frac{\delta^2}{\delta \varphi(x) \, \delta \varphi(y)} \,,$$

 $H_F := \Delta_F + H^{\text{sym}}$ Feynman-like propagator,

Time-Ordered Product. Time Ordering Operator: $\mathcal{T} := \exp(\hbar\Gamma'_{H_r})$.

Causality: $F \cdot_{\mathcal{T}} G = F \star G$ if supp(F) later than ("\geq"\geq") supp(G) $\Rightarrow F \cdot_{\mathcal{T}} G$ uniquely defined iff $supp(F) \cap supp(G) = \emptyset$

Partial Algebra of \mathcal{T} -product: $(\mathcal{F}(\mathbb{M})[[\hbar]], \cdot_{\mathcal{T}})$

pAQFT - Time-Ordered Product

Graph Structure

$$\Gamma_{H_F}(F \otimes G) = \left\langle H_F, F^{(1)} \otimes G^{(1)} \right\rangle = \bullet \bullet, \quad \delta^{\alpha} : \bigotimes_{v \in V} F_v \mapsto \bigotimes_{v \in V} F_v^{(\alpha_v)}$$

Graph expansion

$$F_1 \cdot_{\boldsymbol{\mathcal{T}}} \cdots \cdot_{\boldsymbol{\mathcal{T}}} F_n = \sum_{\alpha \in \mathbb{N}^n} \sum_{\Gamma \in \mathcal{G}_\alpha} \frac{\hbar^{|E(\Gamma)|}}{\operatorname{Sym}(\Gamma)} \left\langle S_\Gamma, \delta^\alpha(F_1 \otimes \cdots \otimes F_n) \right\rangle,$$

 \mathcal{G}_{α} : set of non-tadpole graphs Γ with $|V(\Gamma)| = n = \dim(\alpha)$, $|E(\Gamma)| = \frac{|\alpha|}{2}$

Definition (Local Functional)

 $F \in \mathcal{F}(\mathbb{M})$ is called local $(F \in \mathcal{F}_{loc}(\mathbb{M}))$, if $\forall n \in \mathbb{N}, \forall \varphi \in \mathscr{E}(\mathbb{M})$:

- \circ supp $(F^{(n)}(arphi))\subset \mathsf{Diag}(\mathbb{M}^n)=\{x_1=\cdots=x_n\}$ (thin diagonal)
- WF $(F^{(n)}(\varphi)) \subset [T \operatorname{Diag}(\mathbb{M}^n)]^{\perp}$. e.g., field polynomials.

Implication:
$$F^{(\alpha_v)}(\vec{x}) = \sum_{\vec{k}} f^{\alpha_v, \vec{k}}(x) \, \delta^{(\vec{k})}(\vec{r}) \in \mathscr{D}(\mathbb{M}) \otimes \mathscr{E}'_{\mathsf{Dirac}}(\mathbb{M}^{\alpha_v - 1})$$

Renormalization: Find restriction S_{Γ} of $\bigotimes_{e \in E(\Gamma)} H_F(e)$

Renormalization Group and Main Theorem

Scattering Matrix:

$$S \equiv \exp_{\mathcal{T}} : \mathcal{F}_{loc}(\mathbb{M})[[\hbar]] \to \mathcal{F}(\mathbb{M})[[\hbar]]$$

Perturbative Expansion:

$$\mathcal{S}(F) = \sum_{n=0}^{\infty} \frac{1}{n!} \underbrace{\sum_{\alpha \in \mathbb{N}^n} \sum_{\Gamma \in \mathcal{G}_{\alpha}} \frac{\hbar^{|E(\Gamma)|}}{\mathsf{Sym}(\Gamma)} \left\langle S_{\Gamma}, \delta^{\alpha}(F^{\otimes n}) \right\rangle}_{\mathcal{T}_n(F^{\otimes n})}$$

[C1] Causality.
$$S(A + B) = S(A) \star S(B)$$
 if $supp(A) \gtrsim supp(B)$.

[C2] Starting Element.
$$\mathcal{S}(0)=1$$
, $\mathcal{S}^{(1)}(0)=\mathsf{id}:\mathcal{F}_\mathsf{loc}\to\mathcal{F}_\mathsf{loc}$

[C3]
$$\varphi$$
-Locality. $\mathcal{S}(F)(\varphi_0) = \mathcal{S}(F_{\varphi_0}^{[N]})(\varphi_0) + \mathcal{O}(\hbar^{N+1})$, $F_{\varphi_0}^{[N]}$: Taylor

[C4]
$$\varphi$$
-Independence. $\forall h \in \mathscr{E}(\mathbb{M}) : \left\langle \frac{\delta \mathcal{S}(F)}{\delta \varphi}, h \right\rangle = \mathcal{S}^{(1)}(0)(F) \left\langle \frac{\delta F}{\delta \varphi}, h \right\rangle$

Renormalization Group and Main Theorem Stückelberg-Petermann Group

The Stückelberg-Petermann Renormalization Group: R

$$Z:\mathcal{F}_{\text{loc}}(\mathbb{M})[[\hbar]]\to\mathcal{F}_{\text{loc}}(\mathbb{M})[[\hbar]]\ ,$$

[RG1]
$$Z(0) = 0$$

[RG2] Starting Element. $Z^{(1)}(0) = id$.

[RG3]
$$Z = id + \mathcal{O}(\hbar)$$
.

[RG4] Locality. Let $supp(A) \cap supp(C) = \emptyset$, then

$$Z(A + B + C) = Z(A + B) - Z(B) + Z(B + C)$$
.

[RG5]
$$\varphi$$
-Locality. $Z(F)(\varphi_0) = Z(F_{\varphi_0}^{[N]})(\varphi_0) + \mathcal{O}(\hbar^{N+1})$.

[RG6]
$$\varphi$$
-Independence. $\forall \varphi \in \mathscr{E}(\mathbb{M}): \frac{\delta Z}{\delta \varphi} = 0.$

Renormalization Group and Main Theorem

Stora's Main Theorem

Main Theorem of Perturbative Renormalization

Given two scattering matrices, \mathcal{S} and $\widehat{\mathcal{S}}$, both fulfilling the conditions Causality, Starting Element, φ -Locality, and φ -Independence, [C1]-[C4], there is an element of the Stückelberg-Petermann Renormalization Group, $Z \in \mathcal{R}$, such that

$$\widehat{\mathcal{S}} = \mathcal{S} \circ Z$$
.

Conversely, given a scattering matrix S fulfilling [C1]-[C4] and a $Z \in \mathcal{R}$, then $S \circ Z$ is a scattering matrix fulfilling [C1]-[C4].

Proof: Brunetti, Dütsch, Fredenhagen 2004 / 2009.

Main Theorem in Components

Corollary (Faà di Bruno's Formula)

Let \mathcal{S} fulfill [C1]-[C4], and $Z \in \mathcal{R}$, then

$$(\boldsymbol{\mathcal{S}} \circ \boldsymbol{\mathcal{Z}})^{(n)}(0) = \sum_{\mathcal{P} \in \mathsf{Part}\{1,\dots,n\}} \boldsymbol{\mathcal{S}}^{(|\mathcal{P}|)}(0) \cdot \left(\bigodot_{I \in \mathcal{P}} \boldsymbol{\mathcal{Z}}^{(|I|)}(0) \right),$$

where $Part\{1, ..., n\}$ is the set of all partitions of the index set $\{1, ..., n\}$.

" \odot " denotes symmetrized tensor product: $A \odot B := \frac{1}{2} (A \otimes B + B \otimes A)$.
" \odot " denotes composition of linear maps

$$\mathcal{F}_{\mathsf{loc}}(\mathbb{M})[[\hbar]]^{\otimes n} \xrightarrow{\bigcirc_{I \in \mathcal{P}} Z^{([I])}(0)} \mathcal{F}_{\mathsf{loc}}(\mathbb{M})[[\hbar]]^{\otimes |\mathcal{P}|} \xrightarrow{\mathcal{S}^{([\mathcal{P}])}(0)} \mathcal{F}(\mathbb{M})[[\hbar]].$$

Renormalization Hopf algebra

Commutative Part

Definition (5, commutative part)

Let $\mathfrak H$ be the free commutative algebra generated by differential operators

$$a_n: Z \mapsto a_n(Z) := Z^{(n)}(0),$$

 $\mathsf{product}\colon \left(a_k\odot a_l\right)(Z) := a_k(Z)\odot a_l(Z); \; \mathsf{unit}\colon \; \mathbb{1}(Z) = \mathsf{id} : \mathcal{F}_{\mathsf{loc}} \to \mathcal{F}_{\mathsf{loc}}.$

Coproduct. Faà di Bruno,
$$\Delta a_n := \sum_{\mathcal{P} \in \mathsf{Part}\{1,...,n\}} a_{|\mathcal{P}|} \otimes \bigodot_{I \in \mathcal{P}} a_{|I|}.$$

Grading. $\deg(a_n) := n-1$, compatible with Δ and M_{\odot} . Starting Element condition: $a_1(Z) = Z^{(1)}(0) = \mathrm{id}$ $\Rightarrow (\mathfrak{H}, \odot, \Delta)$ is \mathbb{N}_0 -graded connected $(a_1 = 1)$.

$$\mathsf{Antipode.}\ \ \mathcal{A}(a_n) = -\sum_{\mathcal{P} \in \mathsf{Part}\{1,\dots,n\} \setminus \{\mathcal{P}_1\}} a_{|\mathcal{P}|} \odot \left(\bigodot_{I \in \mathcal{P}} \mathcal{A}(a_{|I|}) \right)$$

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Renormalization Hopf algebra

Noncommutative Part

Composition product. Necessary for dualizing whole structure of \mathcal{R} .

$$\mathtt{C}:\mathfrak{H}\otimes\mathfrak{H} o\mathfrak{H} \qquad , \qquad \mathtt{C}(a_k\otimes igodot_{i=1}^k a_{l_i})=a_k \odot igodot_{i=1}^k a_{l_i}\,,$$

with
$$\left(a_k \odot \bigcirc_{i=1}^k a_{l_i}\right)(Z) := a_k(Z) \cdot \bigcirc_{i=1}^k a_{l_i}(Z)$$
.

Compatibility. with grading, unit, and coproduct.

Convolution. $\phi, \psi \in \operatorname{Aut}(\mathfrak{H}): \phi \bullet_{\mathtt{C}} \psi := \mathtt{C} \circ (\phi \otimes \psi) \circ \Delta.$

Antipode.

$$\mathcal{A}_{\mathtt{C}}(a_n) = -\sum_{\mathcal{P} \in \mathsf{Part}\{1,...,n\} \setminus \{\mathcal{P}_1\}} a_{|\mathcal{P}|} \widehat{\mathbb{C}} \bigodot_{I \in \mathcal{P}} a_{|I|} \,.$$

Algebraic Dual of Stückelberg-Petermann Group: $(\mathfrak{H},\odot,\overline{\mathbb{C}},\Delta,\mathcal{A}_{\mathtt{C}})$

Dimensionally Regularized S-matrix

Choosing a Feynman Propagator

Dimensionally Regularized Feynman Propagator: $d \in 2\mathbb{N}$

$$\begin{split} H_F^{m,\mu,\zeta}(x) &= \underbrace{W_F^{m,\mu,\zeta}(x)}_{\sim K_{\underbrace{d+\zeta}-1}} + \underbrace{B^{m,\mu,\zeta}(x)}_{\stackrel{\sim I_{\underbrace{d+\zeta}-1}}{2}-1} & \in \mathscr{D}'(\mathbb{M}) \,, \quad \zeta \in \Omega \backslash \, \{0\} \subset \mathbb{C} \,. \\ W_F^{m,\mu,\zeta} &\xrightarrow{\zeta \to 0} \Delta_F^{m,\mu}. & B^{m,\mu,\zeta} \colon \text{smooth, symmetric,} \\ &\text{Lorentz invariant Klein-Gordon solution.} \end{split}$$

- $H_F^{m,\mu,\zeta}$ as well as $H_F^{m,\mu,\zeta o 0}$ depend smoothly on mass parameter m^2 .
- $\forall k \in \mathbb{N}: \left[H_F^{m,\mu,\zeta}\right]^k \in \mathscr{D}'(\mathbb{M}\setminus\{0\})$ has unique extension to $\mathscr{D}'(\mathbb{M})$.
- $\forall f \in \mathscr{D}(\mathbb{M}) \colon \zeta \mapsto \left\langle \left[H_F^{m,\mu,\zeta} \right]^k, f \right\rangle \text{ is analytic in } \Omega \backslash \left\{ 0 \right\} \subset \mathbb{C},$ with $\operatorname{pp}\left(\left[H_F^{m,\mu,\zeta} \right]^k \right) \in \mathscr{E}_{\mathsf{Dirac}}'$ a local distribution.

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Dimensionally Regularized S-matrix Scattering Matrix

Regularized Time-Ordering Operator

$$\boldsymbol{\mathcal{T}}^{\zeta} := \exp(\hbar \Gamma'_{H_F^{m,\mu,\zeta}}) \,, \quad \Gamma'_{H_F^{m,\mu,\zeta}} := \frac{1}{2} \int dx \, dy \, H_F^{m,\mu,\zeta}(x,y) \, \frac{\delta^2}{\delta \varphi(x) \, \delta \varphi(y)} \,.$$

Regularized Time-Ordered Product

$$F \cdot_{\mathcal{T}^{\zeta}} G := \mathcal{T}^{\zeta} \left(\mathcal{T}^{\zeta^{-1}} F \cdot \mathcal{T}^{\zeta^{-1}} G \right).$$

Regularized Scattering Matrix

$$\mathcal{S}_{\zeta}(F) := \exp_{\mathcal{T}^{\zeta}}(F)$$
.

- satisfies conditions [C1]-[C4].
- is a regularization outside large diagonal. Let $F_1, \ldots, F_n \in \mathcal{F}_{loc}$ s.t. $\forall i \neq j$: supp $(F_i) \cap \text{supp}(F_i) = \emptyset$, then

$$\lim_{\zeta \to 0} \mathcal{S}_{\zeta}^{(n)}(0)(F_1 \otimes \cdots \otimes F_n) = \mathcal{T}_{\mu}^n(F_1 \otimes \cdots \otimes F_n).$$

Recursion for Counterterms

There is a $Z_{\zeta} \in \mathcal{R}$ such that:

$${\mathcal S}^\mu = \lim_{\zeta o 0} {\mathcal S}_\zeta \circ Z_\zeta$$
 .

Writing this in components, we get

$$(\boldsymbol{\mathcal{S}}_{\zeta} \circ Z_{\zeta})^{(n)}(0) = \sum_{\mathcal{P} \in \mathsf{Part}\{1,\dots,n\}} \boldsymbol{\mathcal{S}}_{\zeta}^{(|\mathcal{P}|)}(0) \cdot \left(\bigodot_{I \in \mathcal{P}} Z_{\zeta}^{(|I|)}(0) \right).$$

Corollary (Recursion for MS counterterms)

The counterterms of minimal subtraction are given by, $\mathcal{P}_1 = \{\{1,\dots,n\}\}$,

$$Z_\zeta^{(n)}(0) = -\mathsf{pp}\left[\sum_{\mathcal{P}\in\mathsf{Part}\{1,\dots,n\}\setminus\{\mathcal{P}_1\}} \boldsymbol{\mathcal{S}}_\zeta^{(|\mathcal{P}|)}(0) \cdot \left(\bigodot_{I\in\mathcal{P}} Z_\zeta^{(|I|)}(0)\right)\right]\,.$$

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Interlude: Forest Formula

The Epstein-Glaser Forest Formula

MS gives *unique* extension at each step of EG, e.g., $Z_{\zeta}^{(2)} = -\text{pp}(\boldsymbol{\mathcal{S}}_{\zeta}^{(2)})$.

 \Rightarrow Solution of Epstein-Glaser recursion is possible.

MS operator $\forall P \in Part\{1, ..., n\}$:

$$-T_{\mathcal{P}}^{\mathsf{MS}}\boldsymbol{\mathcal{S}}_{\zeta}^{(n)} := \boldsymbol{\mathcal{S}}_{\zeta}^{(|\mathcal{P}|)} \cdot \left(\bigodot_{l \in \mathcal{P}} \overline{R}_{|I|} \left(\boldsymbol{\mathcal{S}}_{\zeta}^{(|I|)} \right) \right) \, , \; \overline{R}_{k} = \begin{cases} \mathsf{id} & \mathsf{if} \; k = 1 \\ -\mathsf{pp} & \mathsf{if} \; k > 1 \end{cases}$$

Forest $\mathbb{F}:=\{\cdots\geq\mathcal{P}_n\}$, tot. ord. set of partitions, $\mathcal{P}_n=\{\{1\},\ldots,\{n\}\}$

Theorem (Epstein-Glaser Forest Formula)

The finite part of the n-fold regularized time-ordered product is given by

$$\boldsymbol{\mathcal{S}}_{\zeta,\mathsf{ren}}^{(\textit{n})} := \sum_{\mathbb{F} \subset \mathsf{Part}\{1,\ldots,n\}} \left(\prod_{\mathcal{P} \in \mathbb{F}}^{\geq} - \mathcal{T}_{\mathcal{P}}^{\mathsf{MS}}\right) \boldsymbol{\mathcal{S}}_{\zeta}^{(\textit{n})} \,.$$

Renormalization Hopf algebra $(\mathfrak{H}, \odot, \mathbb{C}, \Delta, \mathcal{A}_{\mathtt{C}})$

Feynman Rules and Minimal Subtraction

- Z-Feynman rules. feyn_Z: $a_n \mapsto a_n(Z) : \mathcal{F}_{loc}^{\otimes n} \to \mathcal{F}_{loc}$, linear.
 - Homomorphism w.r.t. ⊙ and ⓒ

Regularized Feynman rules. feyn $_{\zeta}: a_n \mapsto a_n(\mathcal{S}_{\zeta}): \mathcal{F}_{\text{loc}}^{\otimes n} \mapsto \mathcal{F}$, linear.

Homomorphism w.r.t.

Composition with Renormalization Map.

$$ig(extstyle extstyle (R \circ extstyle (B \circ extstyle extstyle$$

 \circ Homomorphism w.r.t. \odot and \bigcirc \Rightarrow Rota-Baxter arg. redundant.

Counterterms.
$$Z_{\zeta}^{(n)} = \mathcal{A}_{\mathtt{C}}^{\mathtt{feyn}_{\zeta}}(a_n) := R \circ \mathtt{feyn}_{\zeta} \circ \mathcal{A}_{\mathtt{C}}(a_n).$$

Finitely Regularized \mathcal{T} -Product. $\mathcal{S}_{\zeta,\mathrm{ren}}^{(n)} = \mathrm{feyn}_{\zeta} \bullet_{\mathbb{C}} \mathcal{A}_{\mathbb{C}}^{\mathrm{feyn}_{\zeta}}(a_n)$.

Renormalization Hopf algebra

Graph structure

Graphs: $\Gamma \in \mathcal{G}_{\alpha}$: non-tadpole graph; γ_{I} : full vertex part of $I \subset V(\Gamma)$. Part $^{c}V(\Gamma)$: partitions into connected subgraphs, i.e., $\forall \mathcal{P} \in \operatorname{Part}^{c}V(\Gamma): \forall I \in \mathcal{P}: \gamma_{I} \text{ connected}$. Γ/\mathcal{P} : graph with blocks $I \in \mathcal{P}$ as vertices and as lines all lines between different blocks $I, I' \in \mathcal{P}$.

$$\mathsf{Coproduct} \colon \Delta \Gamma = \sum_{\mathcal{P} \in \mathsf{Part}^c V(\Gamma)} \Gamma / \mathcal{P} \otimes \bigcup_{I \in \mathcal{P}} \gamma_I \,.$$

Finitely Regularized Amplitude:

$$\left(\mathtt{feyn}_{\zeta} \, \bullet_{\mathtt{C}} \, \mathcal{A}_{\mathtt{C}}^{\mathtt{feyn}_{\zeta}}\right)(\Gamma) = \sum_{\mathcal{P} \in \mathsf{Part}^{c} \, V(\Gamma)} \mathtt{feyn}_{\zeta}(\Gamma/\mathcal{P}) \cdot \left(\bigodot_{I \in \mathcal{P}} \mathcal{A}_{\mathtt{C}}^{\mathtt{feyn}_{\zeta}}(\gamma_{I}) \right).$$

Example:

Coproduct:

Finitely Regularized Amplitude:

$$\begin{split} \operatorname{feyn}_{\zeta} \bullet_{\operatorname{C}} \mathcal{A}_{\operatorname{C}}^{\operatorname{feyn}_{\zeta}}(\bigodot) &= \operatorname{feyn}_{\zeta}(\bigodot) + \mathcal{A}_{\operatorname{C}}^{\operatorname{feyn}_{\zeta}}(\bigodot) \\ &+ 2\operatorname{feyn}_{\zeta}(\bigodot) \cdot \left(\operatorname{id} \odot \mathcal{A}_{\operatorname{C}}^{\operatorname{feyn}_{\zeta}}(\bigodot) \right) \\ &+ \operatorname{feyn}_{\zeta}(\bigotimes) \cdot \left(\operatorname{id} \odot \mathcal{A}_{\operatorname{C}}^{\operatorname{feyn}_{\zeta}}(\bigodot) \odot \operatorname{id} \right) \end{split}$$

Conclusion

- Hopf algebra structure of pQFT may be understood as the algebraic dual of the Stückelberg-Petermann renormalization group.
- To describe multiple interactions the introduction of an additional, non-commutative Hopf algebra product is necessary.
- Feynman rules and renormalization map emerge naturally and give a recursion for minimal subtraction (MS) counterterms.
- Epstein-Glaser recursion can be solved in terms of a forest formula.
- Connes-Kreimer theory of renormalization follows as a special case from Stora's main theorem.

Outlook

- Applications to the (computer-assisted) computation of higher order contributions to the perturbative expansion?
- Relation to topics in pure mathematics such as Number Theory and Noncommutative Geometry (Multiple Zeta Values, Polylogarithms, Graph Polynomials, ...) as suggested by the Connes-Kreimer framework?
- Renormalization of gauge theories also in the algebraic setting?

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